NEAR-INFRARED SPECTROSCOPY AND A SEARCH FOR CO EMISSION IN THREE EXTREMELY LUMINOUS IRAS SOURCES: IRAS F09105 + 4108, IRAS F15307 + 3252, AND PG 1634 + 706

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ABSTRACT

Rest-frame 0.48–1.1 μm emission-line strengths and molecular gas mass (H2) upper limits for three luminous infrared sources—the hyperluminous infrared galaxies IRAS F09105 + 4108 (z = 0.4417), IRAS F15307 + 3252 (z = 0.2926), and the optically selected QSO PG 1634 + 706 (z = 1.338)—are presented. Diagnostic emission-line ratios ([O III] λ5007/Hβ, [S II] λλ6716, 6731/Hz, [N II] λ6583/Hz, and [S II] λλ9069, 9532/Hz) indicate a Seyfert 2–like spectrum for both infrared galaxies, consistent with previously published work. The upper limits of molecular gas mass for all three sources are \( M_{\text{H}_2} < (1-3) \times 10^7 M_\odot \) (\( q_0 = 0.5, H_\text{o} = 100 \text{ km s}^{-1} \text{ Mpc}^{-1} \)), less than that of the most gas-rich infrared galaxies in the local universe. All three sources have \( L_\text{IR}/L_\text{CO} \sim 1300–2000 \) and thus are the extragalactic sources with the most extreme \( L_\text{IR}/L_\text{CO} \) values measured to date. Given the relatively warm far-infrared colors for all three objects, much of their infrared luminosity may emanate from a relatively modest amount of warm dust (e.g., \( M_\text{dust} \approx 10^5–10^7 \text{ h}^{-2} M_\odot \), \( T_\text{dust} = 200–100 \text{ K} \)) near the AGN. For F09105 + 4108 and F15307 + 3252, the implied circumnuclear covering factor of this dust is \sim 90%, while for PG 1634 + 706 the covering factor is only \sim 35%.

Subject headings: galaxies: active — galaxies: ISM — infrared: galaxies — ISM: molecules — radio lines: galaxies

1. INTRODUCTION

Studies of ultraluminous infrared galaxies (ULIGs) detected by IRAS have focused on revealing the nature of the host galaxies and the source powering the enormous infrared luminosities. Four forms of spectroscopic evidence for both stellar and nonstellar energy sources exist: (1) Studies of ULIGs comparing the fluxes of hydrogen recombination lines with optical and near-infrared-for

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3 The NRAO is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc.

We designate ultraluminous infrared galaxies (ULIGs) as galaxies with \( L_\text{IR} \geq 10^{12} L_\odot \), where \( L_\odot = 4.8 \times 10^{38} \text{ erg s}^{-1} \) and assuming \( H_\odot = 75 \text{ km s}^{-1} \text{ Mpc}^{-1} \), and \( q_0 = 0.5 \) (e.g., Sanders & Mirabel 1996). Additional infrared galaxy definitions used throughout this paper are as follows: luminous infrared galaxies (LIGs) are galaxies with \( L_\text{IR} = (10^{11}–10^{12}) L_\odot \); hyperluminous infrared galaxies (HyLIGs) are galaxies with \( L_\text{IR} \geq 10^{12} L_\odot \). Throughout this paper, we adopt \( H_\odot = 100 \text{ km s}^{-1} \text{ Mpc}^{-1} \) and \( q_0 = 0.5 \).

bined lines have shown that the dominant ionization process appears to differ from galaxy to galaxy. Emission-line diagnostics have been used to classify these galaxies as H II region–like galaxies (i.e., emission lines induced by hot OB stars), low-ionization nuclear emission-line regions galaxies (LINERs), or Seyferts (e.g., Sanders et al. 1988a; Kim et al. 1995; Kim 1995; Veilleux et al. 1995). (2) Spectral polarimetry of ULIGs often reveal highly polarized continuum and broad emission-line spectrum, indicating the presence of a buried active galactic nucleus (AGN; e.g., Kay & Miller 1989; Hines 1991, 1994; Hines & Wills 1993; Jannuzi et al. 1994). (3) Millimeter-wave (CO) spectroscopy of ULIGs has shown them to be rich in molecular gas, indicating the presence of an enormous amount of material for fueling an AGN and star formation (e.g., Young et al. 1984; Sanders & Mirabel 1985; Sanders et al. 1987; Scoville et al. 1986; Sanders, Scoville, & Soifer 1991; Sanders 1991). In addition, morphological studies of ULIGs have revealed that they commonly have tidal distortions and nearby companions indicative of the merger of two or more galaxies (e.g., Sanders et al. 1988a, 1988b). Sanders et al. (1988a, 1988b) have suggested that ULIGs may be the precursor to optically selected quasars.
The recent discovery of the high-redshift HyLIG IRAS 10214+4724 ($z = 2.286$) and the subsequent detection of molecular gas in the source (Brown & Vanden Bout 1991, 1992; Solomon, Downes, & Radford 1992a), as well as evidence that it harbors a buried AGN (e.g., Elston et al. 1994; Jannuzi et al. 1994), have sparked interest in searching for and studying more luminous IRAS galaxies at early epochs of the universe. To date, there are two other published findings of $z > 0.3$ galaxies discovered by IRAS, F09105+4108 ($z = 0.4417$; Kleinmann & Keel 1987; Kleinmann et al. 1988) and F15307+3252 ($z = 0.926$; Cutri et al. 1994), and surveys are underway to search for additional candidates.\(^5\) In addition, a significant fraction of moderate-to-high redshift QSOs were detected by IRAS, making observations that test the infrared galaxy–QSO evolutionary hypothesis at higher redshift (and higher luminosities) feasible. To date, two local QSOs (Mrk 1014 [Sanders, Scoville, & Soifer 1986a] and I Zw 1 [Barvainis, Alloin, & Antonucci 1989]), one moderate-redshift quasar (3C 48 at $z = 0.3695$ [Scoville et al. 1993; Wink, Guilloteau, & Wilson 1997]), and two higher redshift QSOs (the gravitationally lensed H 1413+117 at $z = 2.556$ [Barvainis et al. 1994] and BR 1202–07 at $z = 4.695$ [Ohta et al. 1996; Omont et al. 1996a]) have been unambiguously detected in CO.

In this paper, we present CO observations and near-infrared spectroscopy of the high-redshift HyLIGs F09105+4108, F15307+3252, and PG 1634+706 (IRAS 16347+7037; e.g., Sanders et al. 1989), the latter of which is the most infrared luminous, optically selected QSO detected by IRAS in the PG sample. The purpose of these observations is to extend the determination of molecular gas properties and dominant ionization mechanisms to distant sources with extreme $L_{\text{IR}}$, as well as to establish whether their properties are consistent with the trends in such properties observed for local ULIGs and QSOs.

This paper is divided into five sections. The near-infrared and millimeter observing procedures are discussed in § 2. The data reduction methods and line-flux calculations are summarized in § 3. Section 4 begins with a brief discussion of the optical and infrared morphologies of the three sources; then the dominant ionization mechanism in the two IRAS galaxies, the interpretation of the $L_{\text{CO}}$ and $L_{\text{IR}}$ of the three sources, and their spectral energy distributions are examined. Section 5 summarizes the results. Throughout this paper, we adopt $H_0 = 100$ km s$^{-1}$ Mpc$^{-1}$ and $\Omega_0 = 0.5$.

## 2. Observations

Tables 1–4 summarize basic information and a journal of observations for the three sources discussed in this paper. Observations with each of the telescopes used are discussed separately below. A more in depth discussion of the observing technique and the difficulties inherent in millimeter-wave and submillimeter-wave observations of faint sources can be found in Evans et al. (1996).

### 2.1. Imaging

All of the imaging data were obtained at the UH 2.2 m telescope on Mauna Kea. The near-infrared imaging data were taken with the UH quick infrared camera (QUIRC); Hodapp et al. 1996), which consists of a 1024 $\times$ 1024 pixel HgCdTe astronomical wide area infrared imaging array. The R-band image of F15307+3252 and the I-band image of PG 1634+706 were taken with the UH Tektronix 2048 $\times$ 2048 CCD camera (Wainscoat 1996). Finally, the I-band image of PG 1634+706 was obtained with the Orbit Semicondor 2048 $\times$ 2048 CCD camera (Wainscoat 1996), reimaged to f/10. With the exception of the R-band image of F15307+3252, which consisted of one long exposure, all observations were done by taking a series of dithered exposures.

### 2.2. Near-Infrared Spectroscopy

#### 2.2.1. UH 2.2 m Telescope

All spectroscopic observations on the UH 2.2 m telescope were made with the K-band spectrometer (KSPEC; Hodapp et al. 1994). KSPEC is a cross-dispersed echelle configured to cover the wavelength range 1.1–2.5 $\mu$m in three orders ($J$, $H$, and $K$) on a 256 $\times$ 256 NICMOS-3 HgCdTe detector array. The wavelength resolutions (at 2.2 $\mu$m) for each of the observations are listed in Table 3. Additionally, the wavelength range 0.7–1.0 $\mu$m is also dispersed in several orders on the array. However, the order crowding at 0.7–1.0 $\mu$m is such that only spectra of pointlike sources can effectively be extracted, and this wavelength range is also compromised because of the fact that KSPEC is only in focus at $\lambda > 0.85 \mu$m. In addition to its spectral capabilities, KSPEC also provides simultaneous imaging of approximately 1 arcmin$^2$ of sky around the slit on a second NICMOS-3 array. For the 1994 December observing period, the new UH tip-tilt system (Jim et al. 1998) was also implemented, notably improving the tracking, offsetting, and seeing of the observations and data.

Each observing night began with a series of flat-field exposures in the closed dome with incandescent lights turned on, then off. Once the telescope was guiding with the source in the slit, a 180 s exposure was taken, followed by a 180 s exposure 10° off-source. Each spectrum exposure was accompanied by a 140 s image exposure. For the March observations, the pattern for observing was source-sky-source-sky and so on, but the pattern was changed to source-sky-source-sky-sky-sky-sky-source, and so on, for the December observations to minimize spurious features created by the changing sky conditions. Observations of standard A (V) stars near each source were also made to flux calibrate the source spectrum and to remove telluric lines, and observations of an argon lamp were made to wavelength calibrate the KSPEC data.

#### 2.2.2. United Kingdom Infrared Telescope (UKIRT)

Observations of F15307+3252 were made with the upgraded Coadded Grating Spectrometer (CGS4) on the 3.8 m UKIRT in shared-risk time. CGS4 is a 1–5 $\mu$m twodimensional grating spectrometer with a 256 $\times$ 256 InSb array. Because the H$_\alpha$ + [N II] $\lambda 6583$ and [S II] $\lambda\lambda 6716, 6731$ (hereafter [S II] $\lambda 6724$) emission from the source is redshifted to J-band, the 75 line mm$^{-1}$ grating in second order was used. Sampling was done by moving the detector over 2 pixels in 6 steps in the wavelength direction. For the observations, a 1 pixel wide (122) slit was used.

Spectral data was taken in 120 s exposures, with 20 s of each exposure made at each detector position during the 2 pixel sampling. The slit was then moved 15" in the spatial axis direction of the array for the next 120 s exposure. Thus,
### TABLE 1

**SOURCE PARAMETERS**

<table>
<thead>
<tr>
<th>Name (IRAS)</th>
<th>B1950.0</th>
<th>$f_\nu (\mu Jy)$</th>
<th>References</th>
<th>$L_\nu^* (L_\odot)$</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>R.A.</td>
<td>Decl.</td>
<td>12 $\mu$m</td>
<td>25 $\mu$m</td>
</tr>
<tr>
<td>F09105+4108</td>
<td>09 10 32.84</td>
<td>41 08 53.6</td>
<td>0.13 ± 0.03</td>
<td>0.33 ± 0.03</td>
</tr>
<tr>
<td>F15307+3252</td>
<td>15 30 44.63</td>
<td>32 52 50.6</td>
<td>&lt;0.07</td>
<td>0.08 ± 0.07</td>
</tr>
<tr>
<td>F16348+7037 (PG 1634+706)</td>
<td>16 34 51.0</td>
<td>70 37 37.0</td>
<td>0.06 ± 0.01</td>
<td>0.13 ± 0.01</td>
</tr>
</tbody>
</table>

**Note:** Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds.

*a* Assuming $q_0 = 0.5$ and $H_0 = 100$ km s$^{-1}$ Mpc$^{-1}$.

b *IRAS* Point Source Catalog name in Kleinmann et al. (1988) was P09104 + 4109.

**References**—Positions, fluxes, and redshifts: (1) Becker, White, & Helfand (1995); (2) this paper (IRAS FSC); (3) this paper (IRAS SCANPI); (4) this paper (near-infrared emission lines); (5) Tytler & Fan (1992).
the source appears on the array as a positive and negative spectrum after the first frame is subtracted from the second. The source-sky observation sequence was the same as that used in the 1994 December KSPEC observations. Observations of the standard A star HD 162208 were made before and after the observations of F15307 + 3252. A krypton lamp was used as the wavelength calibration source.

We encountered a problem with CGS4: the source flux of the positive spectrum was more than of the negative spectrum flux, implying that the source was only partially in the slit in the second position. Alternatively, the flux of the positive A star spectrum was less than the negative spectrum flux. The slit also appeared to be tilted relative to the array as evident by the position of the OH sky lines. Because of this, only half the data had enough flux to be usable, and no photometry was possible.

2.3. Submillimeter Spectroscopy

2.3.1. James Clerk Maxwell Telescope (JCMT)

Observations with the 15 m JCMT were made using the 1 mm spectral line receiver (A2), together with the Digital Autocorrelation Spectrometer in wide-band mode (750 MHz bandwidth). Because of excess noise near the edge of the passband, the usable bandwidth was only ~700 MHz, which corresponds to a total velocity coverage of ~900 km s\(^{-1}\) at 230 GHz.

All observations were obtained using a nutating subreflector with a chop rate of ~1.25 Hz. Data were stored as 6 minute scans, and a chopper wheel calibration was performed after every other scan. Pointing was monitored every few hours by observations of standard continuum sources and was estimated to be accurate to within a few arcseconds. During the course of the observations, data were also taken with the bandpass centered at velocity offsets ±40 km s\(^{-1}\) from the velocity corresponding to the redshift of the optical emission lines. Shifting the velocities in such a manner minimizes any ripples inherent in the baseline.

2.3.2. Instituto de Radio Astronomía Milimétrica (IRAM) 30 m Telescope

The IRAM telescope has the ability to perform observations using the 3, 2, and 1 mm receivers simultaneously, making use of the two available 512 MHz filterbanks and two 512 MHz autocorrelators. For F09105 + 4108, the 2 and 1 mm receivers were tuned to the redshifted CO (2 → 1) and CO (3 → 2) transitions, respectively. During the observations, pointing was monitored by observing the planets and standard continuum sources and was typically accurate to within 3\(^\circ\).

2.3.3. National Radio Astronomy Observatory (NRAO) 12 m Telescope

Observations of the redshifted CO (2 → 1) emission in PG 1634 + 706 were made with the NRAO 12 m Telescope configured with a dual polarization SIS spectral-line receiver and two 256 × 2 MHz channel filterbanks. Observations were obtained using a nutating subreflector with a chop

<table>
<thead>
<tr>
<th>Source</th>
<th>Telescope</th>
<th>Detector</th>
<th>Pixel Scale* (arcsec pixel(^{-1}))</th>
<th>Band</th>
<th>Date</th>
<th>Time (minutes)</th>
</tr>
</thead>
<tbody>
<tr>
<td>IRAS F09105+4108......</td>
<td>UH 2.2 m</td>
<td>Tek</td>
<td>0.14</td>
<td>I</td>
<td>1997 Mar</td>
<td>26</td>
</tr>
<tr>
<td>IRAS F15307+3252......</td>
<td>UH 2.2 m</td>
<td>Tek</td>
<td>0.22</td>
<td>K</td>
<td>1996 Apr</td>
<td>30</td>
</tr>
<tr>
<td>PG 1634+706</td>
<td>UH 2.2 m</td>
<td>Orbit</td>
<td>0.14</td>
<td>I</td>
<td>1997 Jun</td>
<td>55</td>
</tr>
</tbody>
</table>

* Imaging at the UH 2.2 m telescope is done with a focal ratio of f/10 or f/31. For the CCD observations, images taken at f/31 are binned 2 × 2 pixels.
rate of \( \sim 1.25 \) Hz. Six minute scans were taken, and a calibration was done every other scan. The pointing was checked every few hours by observing Jupiter and was estimated to be accurate to within a few arcseconds. Because the duration of the observation was equivalent to one source transit, no velocity shifts were made.

3. RESULTS

3.1. Imaging Data

The \( I \)-band data reduction was done with IRAF. The data reduction consisted of flat-fielding the individual images, scaling each image to its median value, and shifting and averaging the images.

The \( K' \)-band data reduction was done similar to the \( I \)-band data reduction, except that after the median level of each individual frame was subtracted, the individual frames were averaged together (without spatial shifting) using a min-max averaging routine in order to create a "sky" image. This sky image was subtracted from the individual frames before they were shifted and averaged.

Because the \( R \)-band image of F15307+3252 consisted of one integration, the image was simply flat-fielded and scaled to its median value. Cosmic rays were then removed from the area around the galaxy. With the exception of the

![Fig. 1.—Broad-band images of the IRAS F09105+4108. (a, b) \( I \)-band images with 56" \( \times \) 56" and \( \sim 143" \( \times \) 143" fields of view, respectively. (c, d) \( K' \)-band images with 56" \( \times \) 56" and 180" \( \times \) 180" fields of view, respectively. For all images, north is up and east is to the left.](image-url)
$K'$-band image of PG 1634 + 706, which shows only the unresolved pointlike QSO, the final reduced versions of all the images are shown in Figures 1–3.

3.2. Near-Infrared Spectra

The KSPEC data reduction was done with IRAF. The sky frames were first subtracted from the source frames. The sky subtracted frames were then averaged and divided by the flat fields. Because there were only a few bad pixels and none were on the spectral area containing the source spectrum, the bad pixels were individually set to zero before the spectral orders were extracted with the APALL package. The extracted spectra were then wavelength calibrated and divided by the standard star spectrum, reduced in the same fashion as the source spectrum, to remove any instrumental effects and atmospheric lines (note that the stellar absorption lines in the standard star were removed using the IRAF task SPLIT before the division was done). Because of the narrow width of the slit, no photometry was determined for any of the KSPEC observations. Finally, the spectrum was multiplied by a Planck blackbody spectrum with the same temperature as the standard star.

Table 5 summarizes the emission-line properties determined for F09105 + 4108 and F15307 + 3252. The emission-line fluxes, $f(\lambda)$, line widths, $\Delta v_{\text{FWHM}}$, and root mean square flux densities, $S_{\text{rms}}$, were determined using the IRAF task SPLIT. The signal-to-noise ratios were determined by using the expression $f(\lambda)/\Delta v_{\text{FWHM}} S_{\text{rms}}$.

![Fig. 2.—Broad-band images of the IRAS F15307 + 3252. (a, b) R-band images with 60" × 60" and ~180" × 180" fields of view, respectively. (c, d) $K'$-band images with 60" × 60" and 180" × 180" fields of view, respectively. For all images, north is up and east is to the left.](image-url)
TABLE 5

<table>
<thead>
<tr>
<th>LINE</th>
<th>INSTRUMENT</th>
<th>( \lambda_{\text{obs}} ) (Å)</th>
<th>( f(\lambda)/f(\text{S II}) )</th>
<th>FWHM* (km s(^{-1}))</th>
<th>S/N</th>
<th>IRAS F09105 + 4108</th>
<th>IRAS F15307 + 3252</th>
</tr>
</thead>
<tbody>
<tr>
<td>H#</td>
<td>KSPEC</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>[O ii] ( \lambda 4959 )</td>
<td>KSPEC</td>
<td>7153.1</td>
<td>2.6 ± 0.8</td>
<td>740</td>
<td>3.2</td>
<td>...</td>
<td>9556.5</td>
</tr>
<tr>
<td>[O ii] ( \lambda 5007 )</td>
<td>KSPEC</td>
<td>7224.6</td>
<td>9.0 ± 1.0</td>
<td>980</td>
<td>8.9</td>
<td>...</td>
<td>9644.6</td>
</tr>
<tr>
<td>[O i] ( \lambda 6300 )</td>
<td>KSPEC</td>
<td>...</td>
<td>&lt;0.44</td>
<td>(1000)</td>
<td>...</td>
<td>...</td>
<td>&lt;1.0</td>
</tr>
<tr>
<td>H\alpha + [N ii]</td>
<td>KSPEC</td>
<td>9463.5( ^b )</td>
<td>6.8 ± 0.2</td>
<td>1710</td>
<td>32</td>
<td>12641.8</td>
<td>9.8 ± 0.5</td>
</tr>
<tr>
<td>[N ii] ( \lambda 6548 )</td>
<td>KSPEC</td>
<td>9441.7 ± 2.9</td>
<td>0.62 ± 0.09</td>
<td>1160 ± 150</td>
<td>...</td>
<td>12618 ± 5</td>
<td>1.4 ± 0.2</td>
</tr>
<tr>
<td>H\beta ( \lambda 6563 )</td>
<td>KSPEC</td>
<td>9460.9 ± 1.9</td>
<td>4.4 ± 0.4</td>
<td>1200 ± 100</td>
<td>...</td>
<td>12636 ± 3</td>
<td>5.2 ± 0.5</td>
</tr>
<tr>
<td>[N ii] ( \lambda 6583 )</td>
<td>KSPEC</td>
<td>9492.2 ± 2.9</td>
<td>1.8 ± 0.3</td>
<td>1160 ± 150</td>
<td>...</td>
<td>12686 ± 5</td>
<td>4.1 ± 0.4</td>
</tr>
<tr>
<td>[S ii] ( \lambda 6724 )</td>
<td>KSPEC</td>
<td>9693.9</td>
<td>1.0 ± 0.2</td>
<td>1120</td>
<td>6.2</td>
<td>12960.3</td>
<td>1.0 ± 0.2</td>
</tr>
<tr>
<td>[S ii] ( \lambda 6966 )</td>
<td>KSPEC</td>
<td>13076.9</td>
<td>0.74 ± 0.08</td>
<td>900</td>
<td>8.7</td>
<td>17467.1</td>
<td>1.1 ± 0.3</td>
</tr>
<tr>
<td>[S ii] ( \lambda 6952 )</td>
<td>KSPEC</td>
<td>13174.3</td>
<td>2.0 ± 0.08</td>
<td>930</td>
<td>20</td>
<td>...</td>
<td>...</td>
</tr>
<tr>
<td>He i ( \lambda 10830 )</td>
<td>KSPEC</td>
<td>15615.5</td>
<td>2.1 ± 0.08</td>
<td>1070</td>
<td>27</td>
<td>20873.8</td>
<td>4.8 ± 0.3</td>
</tr>
<tr>
<td>[O i] ( \lambda 6630 )</td>
<td>CGS4</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>12142.8</td>
<td>&lt;1.6</td>
</tr>
<tr>
<td>H\alpha + [N ii]</td>
<td>CGS4</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>12650.4</td>
<td>7.1 ± 0.7</td>
</tr>
<tr>
<td>[S ii] ( \lambda 6714 )</td>
<td>CGS4</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
<td>12959.6</td>
<td>1.0 ± 0.3</td>
</tr>
</tbody>
</table>

Note—The rotational velocities are adopted values based on the average FWHM of the other, stronger emission lines.

\* Uncorrected for instrument broadening.

\( ^b \) Observed wavelength of H\alpha 6563.

\( ^c \) Observed wavelength of [N ii] \( \lambda 6583 \).

\( ^d \) Emission-line values as determined by model fit.

The CGS4 data were reduced in a same manner as the KSPEC data, except that the flat fielding and masking of bad pixels were done automatically after every observation.

Figure 4 shows the spectra of F09105 + 4108 obtained with KSPEC. Six of the eight emission lines detected are at wavelengths corresponding to a source redshift of \( z = 0.4417 \), identical to the observer-frame optical redshift determined by Kleinmann et al. (1988). The discrepancy in the [O iii] \( \lambda \lambda 4959, 5007 \) doublet redshift (\( z = 0.4429 \)) is due to wavelength calibration inaccuracies (the overlapping of
the last few orders made it difficult to determine which argon line was in which order.

Figure 5 shows the spectra of F15307+3252 obtained with KSPEC and the J-band spectrum obtained with CGS4. The KSPEC spectral lines imply a source redshift of $z = 0.926$, and the CGS4 spectra imply a redshift of $z = 0.927$, both consistent with the redshift of 0.926 determined by Cutri et al. (1994). Note that both F09105+4108 and F15307+3252 show strong forbidden line emission. Further, the Hα + [N II] $\lambda 6548$, 6583 and [S II] $\lambda 6724$ lines widths of F09105+4108 (Table 5) are consistent with those measured from median resolution spectroscopy by Liu, Graham, & Wright (1996). The interpretation of the line emission will be discussed in § 4.2.

To determine the fluxes of the Hα and [N II] $\lambda 6548$, 6583 lines of F09105+4108 and F15307+3252, the line profiles were fitted using the IRAF package SPECTFIT. Initial guesses to the fluxes and line widths were entered into the program, and the line flux ratios and wavelength ratio of [N II] $\lambda 6583$ to [N II] $\lambda 6548$ were fixed to 3 and 1.005345 (i.e., the ratio 6583/6548), respectively. In addition, the line widths of the two [N II] lines were fixed to be equal. Iterations were made until a convincing fit was achieved. The best-fit models are listed in Table 5.

Figure 6 is a short exposure of the redshifted [O III] $\lambda\lambda 4959, 5007$ and Hα lines of PG 1634+706. The redshifts determined from the $\Delta F_{\text{FWHM}} = 4600$ km s$^{-1}$ Hα line and the [O III] $\lambda 5007$ lines are $z = 1.339$ and 1.337, respectively. The systemic redshift determined by Tytler & Fan (1992) is 1.3371.

### 3.3. CO Spectra

Table 6 summarizes the CO line data for the three sources. In the case of the IRAM 30 m observations, where data were taken at several widely spaced velocity offsets, the data are presented in separate spectral blocks. Separating the data in this manner avoids the creation of direct current induced ghost features (see Evans et al. 1996) and provides data blocks with relatively constant signal-to-noise across each spectrum.

Figures 7–9 show the spectra of F09105+4108, F15307+3252, and PG 1634+706, respectively. For each source, individual scans were examined to check for sinusoids in the baselines and other irregularities. After bad scans were removed, the rest of the scans were averaged and linear baselines were subtracted. For the JCMT 15 m data, where the sources were observed with several small velocity offsets, only the overlapping regions of the scans were averaged. Finally, the spectra were smoothed to 30–50 km s$^{-1}$. No emission lines were detected. More sensitive upper limits of the CO emission in PG 1634+706 and F15307+3252 have also been recently reported by Barvainis et al. (1998) and Yun, Scoville, & Evans (1998), respectively (see Table 6).

#### 3.3.1. CO Luminescence and H$_2$ Mass Limits

Table 6 lists the root mean square temperatures, $T_{\text{rms}}$, achieved with a velocity resolution, $\Delta v_{\text{res}}$, for the three sources observed. To calculate the upper limit of the CO line intensity, $I_{\text{CO}}$, we assume that the CO lines have a FWHM velocity, $\Delta v_{\text{FWHM}} \sim 250$ km s$^{-1}$, similar to CO
lines seen in local ULIGs (e.g., Sanders et al. 1991) and in IRAS 10214 + 4724 (e.g., Solomon et al. 1992a). Using the observed $T_{\text{rms}}$ and $\Delta v_{\text{res}}$ from Table 6, the 3 $\sigma$ limit on $I_{\text{CO}}$ is given by

$$I_{\text{CO}} < \frac{3T_{\text{rms}} \Delta v_{\text{FWHM}}}{\sqrt{\Delta v_{\text{FWHM}}/\Delta v_{\text{res}}}} \text{ K km s}^{-1}.$$  (1)

Multiplying $I_{\text{CO}}$ by the kelvin-to-jansky conversion for a point source$^6$ gives the corresponding limit for $S_{\text{CO}} \Delta v$ Jy km s$^{-1}$.

$^6$ Kelvin-to-jansky conversion factors: JCMT 16 m, 156 Jy K$^{-1}$ (Matthews 1992); IRAM 30 m, 4.6 to 4.9 Jy K$^{-1}$ (Guelin, Kramer, & Wild 1995); NRAO 12 m, 25.2 Jy K$^{-1}$ (P. Jewell 1995, private communication).

The luminosity distance for a source at a given redshift, $z$, is

$$D_L = cH_0^{-1}q_0^{-2} \times [zq_0 + (q_0 - 1)(\sqrt{2q_0 z + 1} - 1)] h^{-1} \text{ Mpc},$$  (2)

where $q_0$ is the deceleration parameter ($= 0.5$ for a critical density universe) and $H_0 = 100\ h$ (km s$^{-1}$ Mpc$^{-1}$). Given the estimated upper limit of $S_{\text{CO}} \Delta v$, we can calculate the CO luminosity of a source at redshift $z$:

$$L_{\text{CO}} = \left(\frac{c^2}{2kv_{\text{obs}}^2}\right) S_{\text{CO}} \Delta v D_L^2 (1 + z)^{-3},$$  (3)
TABLE 6

CO EMISSION-LINE DATA

<table>
<thead>
<tr>
<th>Source</th>
<th>Site</th>
<th>Line</th>
<th>$v_{\text{rest}}$ (GHz)</th>
<th>$v_{\text{obs}} - v_{\text{rest}}$ (km s$^{-1}$)</th>
<th>$\Delta v_{\text{FWHM}}$ (km s$^{-1}$)</th>
<th>$T_{\text{mb}}$ (mK)</th>
<th>$S_{\text{co}}$ ($\times 10^{20}$ h$^{-2}$ Mpc$^{-2}$)</th>
<th>$L'<em>{\text{CO}}$ ($\times 10^{6}$ h$^{-2}$ M$</em>{\odot}$)</th>
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</thead>
<tbody>
<tr>
<td>IRAS F09105+4018......</td>
<td>15 m</td>
<td>CO ($2 \rightarrow 1$)</td>
<td>239.853</td>
<td>-104</td>
<td>37</td>
<td>1.7</td>
<td>&lt;7.7</td>
<td>&lt;3.0</td>
</tr>
<tr>
<td></td>
<td>30 m</td>
<td>CO ($2 \rightarrow 2$)</td>
<td>239.853</td>
<td>0</td>
<td>31</td>
<td>6.2</td>
<td>&lt;8.2</td>
<td>&lt;3.1</td>
</tr>
<tr>
<td></td>
<td>30 m</td>
<td>CO ($2 \rightarrow 3$)</td>
<td>239.853</td>
<td>625</td>
<td>30</td>
<td>3.3</td>
<td>&lt;4.3</td>
<td>&lt;1.6</td>
</tr>
<tr>
<td>IRAS F15307+3252......</td>
<td>15 m</td>
<td>CO ($4 \rightarrow 3$)</td>
<td>239.415</td>
<td>-108</td>
<td>38</td>
<td>1.2</td>
<td>&lt;5.6</td>
<td>&lt;4.5</td>
</tr>
<tr>
<td></td>
<td>30 m</td>
<td>CO ($4 \rightarrow 3$)</td>
<td>239.415</td>
<td>0</td>
<td>30</td>
<td>5.7</td>
<td>&lt;7.4</td>
<td>&lt;6.0</td>
</tr>
<tr>
<td></td>
<td>188</td>
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<td>8.4</td>
<td>&lt;11</td>
<td>&lt;8.9</td>
<td>&lt;3.5</td>
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<tr>
<td></td>
<td>-188</td>
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<td>8.4</td>
<td>&lt;11</td>
<td>&lt;8.8</td>
<td>&lt;3.5</td>
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<tr>
<td></td>
<td>814</td>
<td>30</td>
<td>3.6</td>
<td>&lt;4.7</td>
<td>&lt;3.8</td>
<td>&lt;1.5</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td>-814</td>
<td>30</td>
<td>3.0</td>
<td>&lt;3.9</td>
<td>&lt;3.1</td>
<td>&lt;1.2</td>
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<td></td>
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<tr>
<td>PG 1634+706 .............</td>
<td>12 m</td>
<td>CO ($2 \rightarrow 1$)</td>
<td>98.647</td>
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<td>49</td>
<td>0.46</td>
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<td>&lt;23</td>
</tr>
<tr>
<td></td>
<td>15 m</td>
<td>CO ($5 \rightarrow 4$)</td>
<td>246.584</td>
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<td>36</td>
<td>1.6</td>
<td>&lt;7.4</td>
<td>&lt;7.1</td>
</tr>
</tbody>
</table>

$^a$ The difference between the velocity the receivers were tuned to and the systemic velocity of the source.

$^b$ The rms temperatures are given in terms of main-beam brightness temperature.

$^c$ Assuming $n_{e} = 0.5$, $H_\alpha = 100 \, \text{km s}^{-1} \, \text{Mpc}^{-1}$.

$^d$ Assuming $x = 4 M_\odot (\text{K km s}^{-1} \text{pc}^{-2})^{-1}$ (see text).

$^e$ The $3 \sigma$ upper limit of $M(H_2)$ reported by Yun et al. (1998) over the redshift range 0.9240–0.9275 is $9.8 \times 10^6 \, h^{-2} \, M_\odot$.

$^f$ The $3 \sigma$ upper limit of $M(H_2)$ reported by Barvainis et al. (1998) is $2.0 \times 10^6 \, h^{-2} \, M_\odot$.

The upper limits on H$_2$ masses implied by the observed limits on $L'_{\text{CO}}$ can be determined by making the reasonable assumption that the CO emission is optically thick and thermalized, and that it originates in gravitationally bound molecular clouds. For molecular gas in gravitationally bound clouds, the ratio of the H$_2$ mass and the CO luminosity is given by $\alpha = M(H_2)/L'_{\text{CO}} \propto [n(H_2)]^{1/2}/T_\text{rot} \times M_\odot (\text{K km s}^{-1} \, \text{pc}^{-2})^{-1}$, where $n(H_2)$ and $T_\text{rot}$ are the density and brightness temperature for the appropriate CO transition (e.g., Scoville & Sanders 1987; Solomon et al. 1992a; Solomon, Radford, & Downes 1992b). Multitransition CO surveys of molecular clouds in the Milky Way (e.g., Sanders et al. 1993) and in nearby starburst galaxies (e.g., Güsten et al. 1993) have shown that hotter clouds tend to be denser in such a way that the density and temperature dependencies tend to cancel each other. The variation in the value of $\alpha$ is less than a factor of 2 for a wide range of gas kinetic temperature, gas densities, and CO abundances. We adopt a value of $4 M_\odot (\text{K km s}^{-1} \, \text{pc}^{-2})^{-1}$ for $\alpha$, which is similar to

![Fig. 6—KSPEC spectrum of [O ii] λ5007 and Hα emission from PG 1634 + 706. The pixel sampling for the H and K′-bands are 12 and 14 μm pixel$^{-1}$, respectively. Note that the broad Fe ii complex is also evident redward of 12000 Å.](image-url)
Fig. 7.—JCMT 15 m and IRAM 30 m spectra of IRAS F09105 + 4108 (z = 0.4417). The intensity scale is main-beam brightness temperature (IRAM) and antenna temperature corrected for aperture losses (JCMT). A linear baseline has been subtracted from each spectrum. The zero velocity corresponds to the average redshift of optical/near-infrared emission lines observed for IRAS F09105 + 4108.

the value determined for the bulk of the molecular gas in the disk of the Milky Way (see Scoville & Sanders 1987), as well as the value determined from a multitransition CO analysis of IRAS 10214 + 4724 (Solomon et al. 1992a). Table 6 lists the upper limits on the molecular gas mass of F09105 + 4108, F15307 + 3252, and PG 1634 + 706 determined using this value of z.

4. DISCUSSION

As stated in the Introduction, spectroscopy has been a powerful tool for determining the dominant ionization source in infrared-luminous galaxies. The environment and morphology of the host galaxy are an essential supplement to the spectroscopy, often providing information on global spatial scales that have direct bearing on the nuclear activity. We briefly discuss the imaging data, then dedicate the bulk of the remaining discussion to the spectroscopic data.

4.1. Morphologies and Environment

Figures 1–3 show images of the three sources taken at
wavelengths longward of $V$-band. Most local ($z < 0.3$) ultraluminous infrared galaxies are observed to have distorted morphologies, multiple nuclei, and close companions indicative of the merger of two or more galaxies (e.g., Sanders & Mirabel 1996 and references therein; Surace et al. 1998). F09105 + 4108 clearly shows evidence for companions (Figure 1). Indeed, unlike most luminous infrared galaxies (LIGs), F09105 + 4108 is clearly a cD galaxy. Both the $I$- and $K$-band images of the galaxy show evidence of “clumps” within the envelope of the cD galaxy, which may be galaxies stripped by tidal or dynamical friction. Such stripping may be the source of dust for F09105 + 4108. In addition, F09105 + 4108 is known to have radio jets (Hines & Wills 1993), indicating the presence of a buried AGN.

While the $R$-band image of F15307 + 3252 (Fig. 2a) shows no direct evidence of distorted morphologies or close companions, the $K$-band image (Fig. 2b) shows a faint companion to the southwest of the main galaxy. Soifer et al. (1996) and Liu et al. (1996) also present near-infrared images showing this companion, with the latter group showing evidence for a second companion. No evidence is seen for tidal features in Figure 2, or in any other images published to date.

The QSO PG 1634 + 706 (Fig. 3), the most distant of the three sources, shows no evidence of distortions or close companions. There appear to be approximately 30 faint objects in the vicinity of the QSO with $I$-band magnitudes, $m_I = 23.5–25.5$, which is consistent with the density of faint
field objects expected within the given magnitude range (Pozzetti et al. 1998). Thus, there is also no apparent evidence with the present data set that the QSO is imbedded in a cluster.

4.2. Emission-Line Diagnostics

Figure 10 shows four emission-line diagnostic diagrams commonly used to distinguish galaxies with H ii region-like properties from Seyfert/LINER-like properties. The usefulness of these diagrams can be understood by the following emission-line characteristics: Emission from [S II] $\lambda 6724$ and [N II] $\lambda 6548$, 6583 can emanate from ionized hydrogen regions, and [S II] $\lambda 6724$ and [O I] $\lambda 6300$ emission can emanate from semi-ionized regions where collisional ionization is significant. Enhancement of these lines occurs in AGNs because, unlike H ii regions, they have extended partially ionized zones created by an excess of X-rays (the absorption cross sections of neutral hydrogen, helium, and all ions are small for X-rays, thus X-rays tend to escape the ionized region before interacting: Veilleux & Osterbrock 1987). [O III] $\lambda 5007$ is a high-ionization line photoionized by UV photons and thus tends to be strong in Seyfert galaxies. [S II] $\lambda 6969$, 9532 is not as sensitive to the level of ionization as [O III] $\lambda 5007$, most likely because of its lower ionization potential (Osterbrock, Tran, & Veilleux 1992). However, [S II] $\lambda 6969$, 9532 and [S II] $\lambda 6724$ together do provide some degree of separation between the three classes of active galaxies (Fig. 10d).

Emission-line ratios of F09105 + 4108 and F15307 + 3252 are plotted in the diagnostic diagrams. For F09105 + 4108, the [O III] $\lambda 5007$/$H\beta$ and [O I] $\lambda 6300$/$H\alpha$ ratios have been determined from Kleinmann et al. (1988). Because their [O III] $\lambda 5007$/$H\beta$ ratio is being used in combination with our KSPEC measurements, it is worth cautioning that [O III] $\lambda 5007$/$H\beta$ may vary as a function of the aperture size used to observe F09105 + 4108. In all four diagrams, both galaxies lie in the Seyfert region, indicating that the dominant source of ionization is a hard, nonthermal continuum source. The analyses of these data are consistent with the detection of broad emission lines in polarized light in all three sources (Hines & Wills 1993; Hines et al. 1995; Jannuzi et al. 1994) and previous optical emission-line diagnostics of both galaxies done with low-resolution (Soifer et al. 1994, 1996) and median-resolution (Kleinmann et al. 1988; Liu et al. 1996) spectroscopy.

Table 7 is a comparison of the observed [O III] $\lambda 5007$/Hz, [S II] $\lambda 6724$/Hz, and [N II] $\lambda 6548$, 6583/$H\alpha$ ratios of F09105 + 4108 and F15307 + 3252 with the sample of bright IRAS galaxies and warm IRAS galaxies (which have been divided into Seyfert 2, LINER, and H ii region classifications). For F09105 + 4108, the ratio [O III] $\lambda 5007$/Hz has been determined from Kleinmann et al. (1988), who

<table>
<thead>
<tr>
<th>Source</th>
<th>$\lambda 5007$/$H\alpha$</th>
<th>$\lambda 6548 + 6583$/$H\alpha$</th>
<th>$\lambda 6724$/$H\alpha$</th>
</tr>
</thead>
<tbody>
<tr>
<td>IRAS F09105 + 4108</td>
<td>3.9 $^*$</td>
<td>0.55 $\pm$ 0.2</td>
<td>0.23 $\pm$ 0.11</td>
</tr>
<tr>
<td>IRAS F15307 + 3252</td>
<td>3.0 $\pm$ 0.4</td>
<td>$&gt; 11.1$ $\pm$ 0.3</td>
<td>1.1 $\pm$ 0.2</td>
</tr>
</tbody>
</table>

$^*$ Determined from [O III] $\lambda 5007$/$H\beta$ = 12 (Kleinmann et al. 1988) and $H\alpha$/$H\beta$ $\sim$ 3.1. The $H\alpha$/$H\beta$ ratio is based on the determination by Kleinmann et al (1988) that $E(B-V) < 0.15$ for F09105 + 4108.
calculate [O iii] λ5007/Hβ ~ 12 and a negligible extinction (note that the [O iii] λ5007/Hβ ratio calculated from our data indicate that half of the [O iii] λ5007 flux is lost via defocusing; see § 22.1). Despite the obvious scatter in the [O iii] λ5007/Hβ ratio due to dust, there is a notable separation between the ratio for Seyferts and for LINERs and H II region galaxies (i.e., the separation is due mainly to the ionization parameter). While F09105 + 4108 has a [O iii] λ5007/Hβ ratio 2 standard deviations above the average ratio for low-redshift Seyfert 2 galaxies, F15307 + 3252 has an observed [O iii] λ5007/Hβ ratio higher than the average Seyfert 2 ratio, but still within the standard deviation of the Seyfert 2 ratios. Inversely, the [S ii] λ6724/Hz and [N ii] λ6548, 6583/Hz ratios are lower than the average values of the Seyferts, but still within the standard deviation of the Seyfert 2 ratios. This may imply that the hyperluminous galaxies have smaller semi-ionized regions or lower metallicities than the average, local IRAS Seyferts.

4.3. $L_\text{IR}$ Versus $L_{\text{CO}}$

Given the evidence presented here and elsewhere for AGNs in both F09105 + 4108 and F15307 + 3252, we consider the infrared and CO luminosities of these two sources and PG 1634 + 706 relative to LIGs, ULIGs, and QSOs detected by IRAS, but having lower infrared luminosities. Figure 11a is a plot of $L_{\text{CO}}$ versus $L_\text{IR}$; the upper limits on the $L_{\text{IR}}/L_{\text{CO}}$ of F09105 + 4108, F15307 + 3252, and PG 1634 + 706 are less than the $L_{\text{CO}}$ of the most gas-rich galaxies observed locally, perhaps indicating that the upward trend of data points in the diagram levels off at higher $L_\text{IR}$ or disperses. Figure 11b is a plot of $L_\text{IR}/L_{\text{CO}}$ versus $L_\text{IR}$; the lower limits on the $L_{\text{IR}}/L_{\text{CO}}$ of the three sources are consistent with the increasing values of $L_\text{IR}/L_{\text{CO}}$ as a function of increasing $L_\text{IR}$ of the sample of LIGs and ULIGs. One possible explanation for the data trends observed in Figure 11 is that $L_\text{IR}$ becomes increasingly dominated by reprocessed nonstellar light at high $L_\text{IR}$. From the spectral energy distributions plotted in Figure 12, it is clear that all three sources have $vL_\nu(60 \mu\text{m})/vL_\nu(100 \mu\text{m})$ ratios higher than nearby, less luminous ULIGs that have been detected in CO, such as Arp 220 and Mrk 1014 (e.g., Sanders et al. 1988b). The mass of “warm” dust required to produce the observed infrared luminosity can be computed by fitting the rest-frame far-infrared (40–
100 μm) emission with a single-temperature dust model using dust emissivities from Draine & Lee (1984). The dust temperatures implied by the far-infrared spectral energy distributions (SEDs) in Figure 12 are in the range 100–200 K, and the corresponding dust masses, \( M_d = (L_d/10^8 L_\odot)(40 \text{ K}/T_d) \), are in the range 10^{-7}–10^{-5} h^{-2} M_\odot. Assuming, as is found for ULIGs like Arp 220 and Mrk 1014, that the bulk of the dust is heated to the computed “warm” dust temperature and that the gas/dust ratio is similar to what is found for other ULIGs such as Arp 220 (i.e., \( \sim 200 \); Sanders et al., 1991), then the implied range of H_2 gas mass is 10^{5}–10^{7} h^{-2} M_\odot. Thus it is possible that the observed infrared luminosities are produced by a relatively modest mass of gas and dust (e.g., comparable to or less than the total molecular gas mass of the Milky Way) heated to temperatures 100–200 K, perhaps by the central AGN. Such small amounts of molecular gas would also explain our nondetections of all three sources in CO (see Table 6). Further, while it is clear from the observed narrow emission lines and large infrared excesses \( (L_{\text{IR}}/L_{\text{bol}} > 0.9) \) of F09105+4108 and F15307+3252 that the circumnuclear covering factor of the dust is almost unity, the broad emission lines and smaller infrared excess \( (L_{\text{IR}}/L_{\text{IR}} < 0.35) \) of PG 1634+706 indicate that the QSO has a low covering factor.

It is curious that CO surveys of galaxies at cosmological distances have yielded no evidence of galaxies with \( L_{\text{CO}} \) much larger than that of the most CO luminous galaxies at lower redshift (\( z < 0.2 \)) plotted in Figure 11 [i.e., \( L_{\text{CO}} \) corresponding to \( (3–4) \times 10^{10} h^{-2} M_\odot \) of molecular gas]. To date, there exists little evidence that the H_2 content of active galaxies evolves with redshift (\( z < 4 \)). The extremely high H_2 masses originally computed for IRAS 10214+4724 and H1413+117 were in large part due to gravitational lensing of these sources by foreground objects (IRAS 10214+4724 [Elston et al. 1994; Soifer et al. 1995; Trentham 1995; Graham & Liu 1995; Broadhurst & Lehar 1995; Serjeant et al. 1995; Close et al. 1995; Eisenhardt et al. 1996; Downes, Solomon, & Radford 1995] and H1413+117 [Barvainis et al. 1994]), and CO surveys of 1 < z < 4 high-redshift powerful radio galaxies (HzPRGs [Evans et al. 1996; Downes et al. 1996]) and z ~ 4 luminous high-redshift QSOs (Barvainis & Antonucci 1996) were unsuccessful at detecting sources with more molecular gas than the most gas-rich galaxies observed locally [i.e., based on 3 \( \sigma \) upper \( M(H_2) \) limits of the HzPRGs and high-redshift QSOs]. The reported detection of CO emission in the z = 4.7 QSO BR 1202–07 \( [M(H_2) \sim 6 \times 10^{10} h^{-2} M_\odot] \) may be the first indication that \( M(H_2) \) increases beyond \( z \sim 4 \). However, reports of shear due to gravitational lensing in the vicinity of BR 1202–07 may mean that the CO luminosity is being amplified (e.g., see the discussion in Omont et al. 1996b). Further, there have been recent, unconfirmed detections of CO emission in the z = 2.4 radio galaxy 53W002 (Scoville et al. 1997), the z = 4.5 QSO BR 1335–07 (Guilloteau et al. 1997), and the gravitationally lensed, z = 2.6 quasar MG 0414+0534 (Barvainis et al. 1998), but all have molecular...
gas masses comparable to the most gas-rich, local galaxies. The inferred molecular gas masses and upper limits of high-redshift galaxies surveyed to date may be evidence that either galaxies with more than \( \sim (3-4) \times 10^{10} h^{-2} M_\odot \) of molecular gas rapidly turn the gas into stars and/or that galaxies form in a hierarchical fashion. It may also be the case that only high-redshift galaxies with a significant nonstellar contribution to their infrared luminosity will be observed to have high \( L_\text{IR} \). A better understanding of the relationship between extreme \( L_\text{IR} \) and \( H_2 \) mass awaits the detection and observation of a larger sample of HyLiGs, especially those having relatively cool IRAS colors, i.e., relatively low values of \( vL_\nu(60 \mu \text{m})/vL_\nu(100 \mu \text{m}) \).

5. SUMMARY

We have presented new millimeter wave CO \((J + 1, J)\) observations and rest-frame 0.48–1.1 \( \mu \text{m} \) spectroscopy of the HyLiGs IRAS F09105 +4108, IRAS F15307 +3252, and PG 1634 +706. The following conclusions are drawn:

1. The emission-line ratios \([ \text{O ii}] \lambda 5007/\text{H} \beta, [ \text{O i}] \lambda 6300/\text{Hz}, [ \text{N ii}] \lambda 6583/\text{Hz}, [ \text{S ii}] \lambda 6724/\text{Hz}, \) and \([ \text{S iii}] \lambda 2300/\text{Hz} \) are correspondingly to a molecular gas mass, \( M_{\text{H}_2} \approx 10^{5.3} - 10^{6.7} h^{-2} M_\odot \). This would be consistent with our measured upper limits for \( L_\text{IR} \) and the corresponding extreme values of \( L_\nu/L_\text{IR} \).

2. While the narrow emission lines and \( L_\nu/L_\text{bol} > 0.90 \) of F09105 +4108 and F15307 +3252 indicate a circumstellar dust covering factor of almost unity, the broad emission lines and \( L_\nu/L_\text{bol} \sim 0.35 \) of PG 1634 +706 indicate that most of the AGN is exposed.
IRAS NEAR-IR SPECTROSCOPY AND CO EMISSION